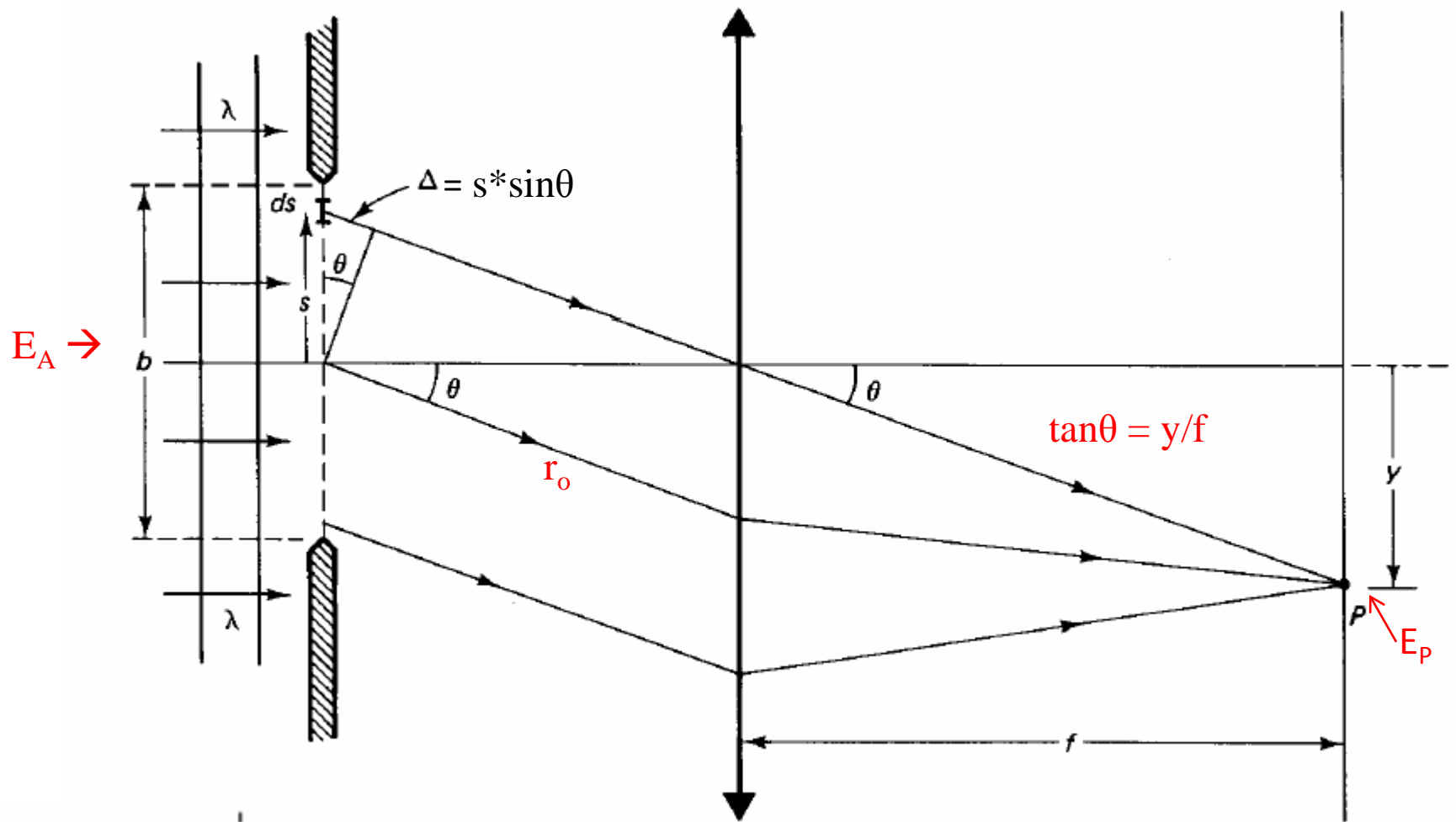


Motion of Diffraction Pattern on VIIRS Detectors

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$$E_P = (E_A/r_o) * \iint \exp(ik * \Delta) dA = (E_A/r_o) * \iint \exp(ik * (s * \sin \theta)) dA$$

$$dA = x ds = 2 * \sqrt{R^2 - s^2} ds$$

Let $v = s/R$, $\gamma = kR \sin \theta$, then

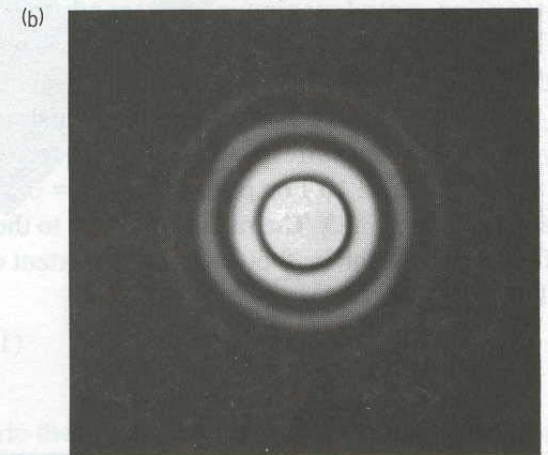
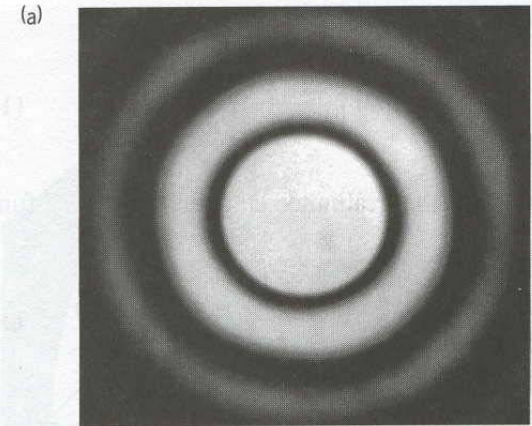
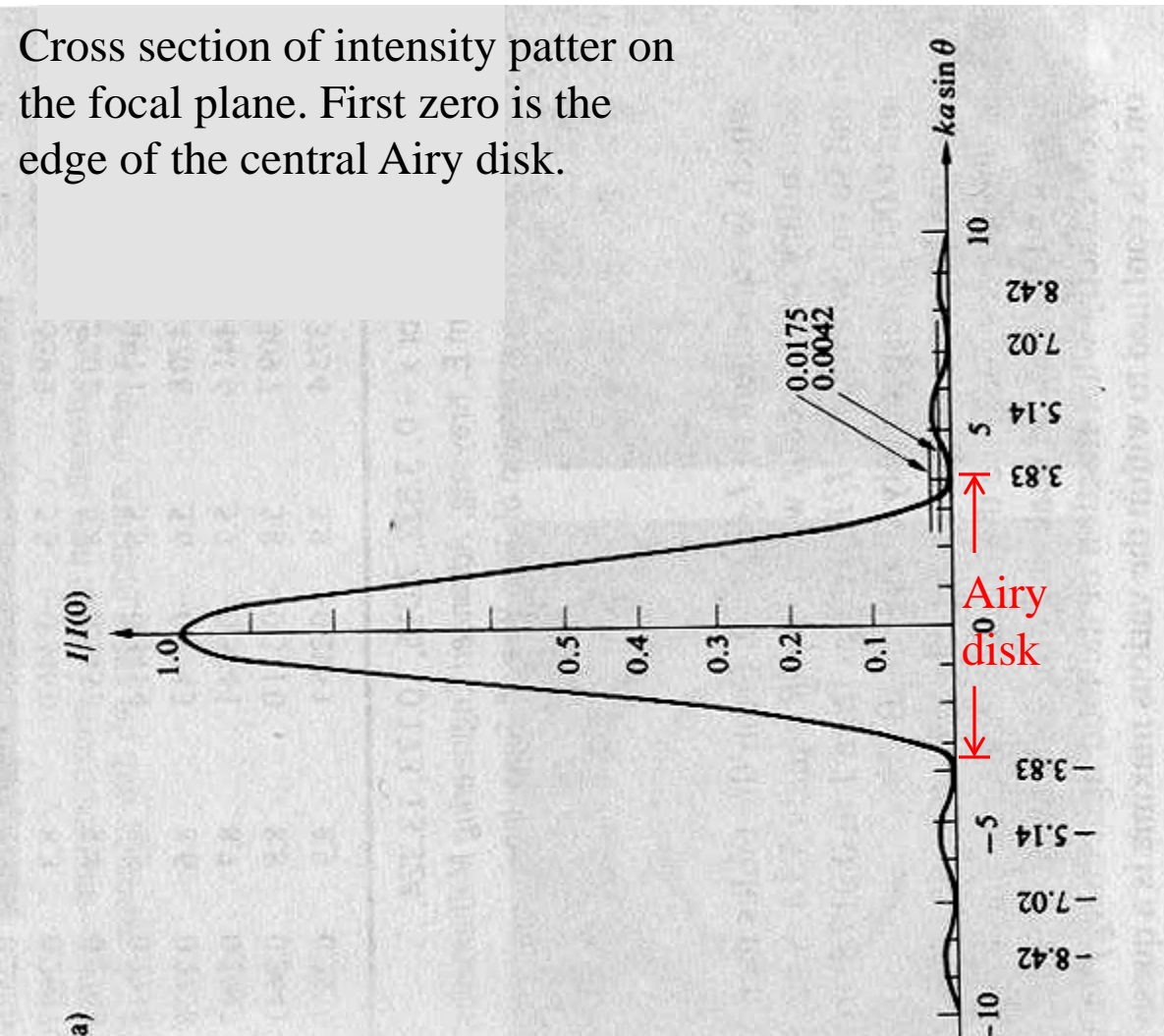
$$E_P = 2E_A R^2 / r_o \left(\int_{-1}^1 \exp(i\gamma v) * \sqrt{1 - v^2} dv \right)$$

$$= 2E_A R^2 / r_o \left(\pi J_1(\gamma) / \gamma \right), \text{ thus}$$

$$I(\theta) = I(0) * \left(2J_1(\gamma) / \gamma \right)^2$$

Plan view of intensity pattern.
Inner most dark circle bounds
the interior bright Airy disk.

Cross section of intensity pattern on
the focal plane. First zero is the
edge of the central Airy disk.



!

$dx = 50.0$! Meters

$distance = dx * \sqrt{dfloat((i-1)-(center-1))**2.0 + dfloat((j-1)-(center-1))**2.0}$

!

!*** From slide 2 we have the following:

!

$\theta = \text{atan}(distance/h)$ (On Earth). $\theta = \text{atan}(y/f)$ (On focal plane)

$\sin_theta = \sin(\theta)$

$\gamma = kR\sin\theta = 2\pi/\lambda * R\sin\theta = \pi d/\lambda * \sin\theta$ Aside: $\theta_R \sim \sin\theta_R = \gamma_1 \lambda/\pi d \sim 3.83 \lambda/\pi d \sim 1.22 \lambda/d$

!

Where θ_R is radius of Airy disk; γ_1 first zero of J1

! *** We'll replace γ with “phase” below.

!

$phase = ((\pi*d)/(wavelength)) * \text{abs}(\sin_theta)$

call Int_Bessel(m,phase,Y)

$J1(i,j) = Y(1)$

$Intensity(i,j) = (2.0*Y(1)/phase)**2.0$

For calculations: $\lambda = 3.9e-6$ (m) (wavelength of incident radiation—*on detector in satellite*)

$h = 35786e3$ (m) (distance from satellite to the surface of the Earth)

$d = 0.3048$ (m) (12 inch aperture of telescope)

$dx = 50.0$ (m) (length—*on the surface of the Earth*)

Diffraction patterns for GOES-16 ABI

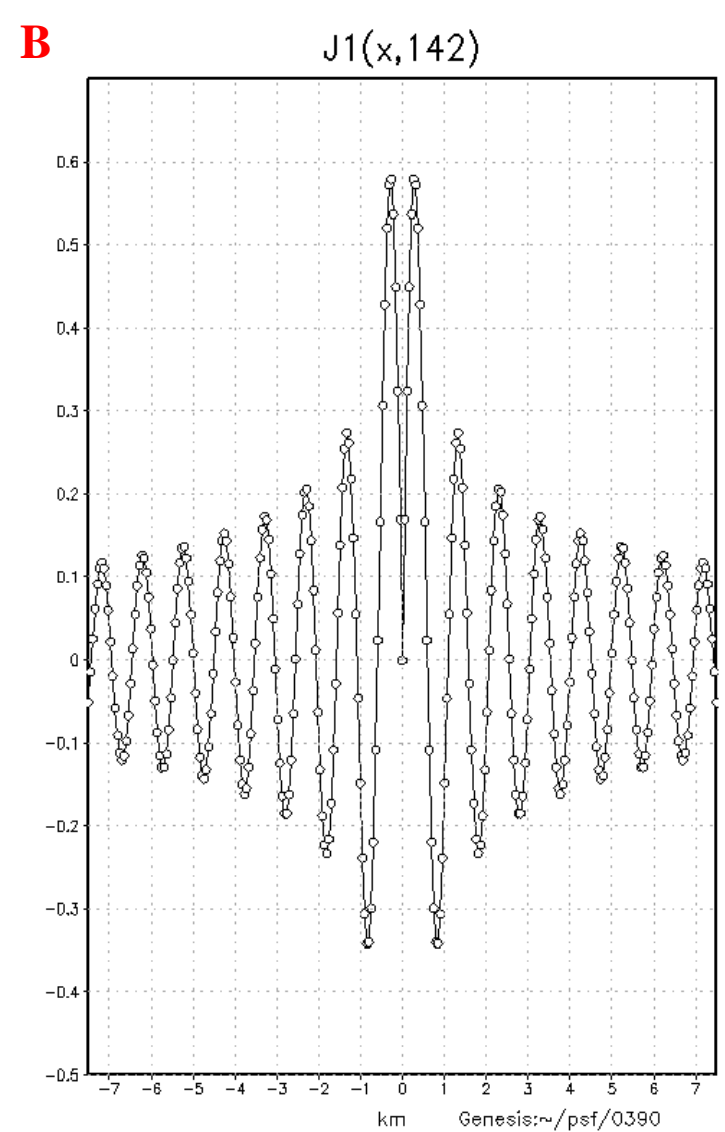
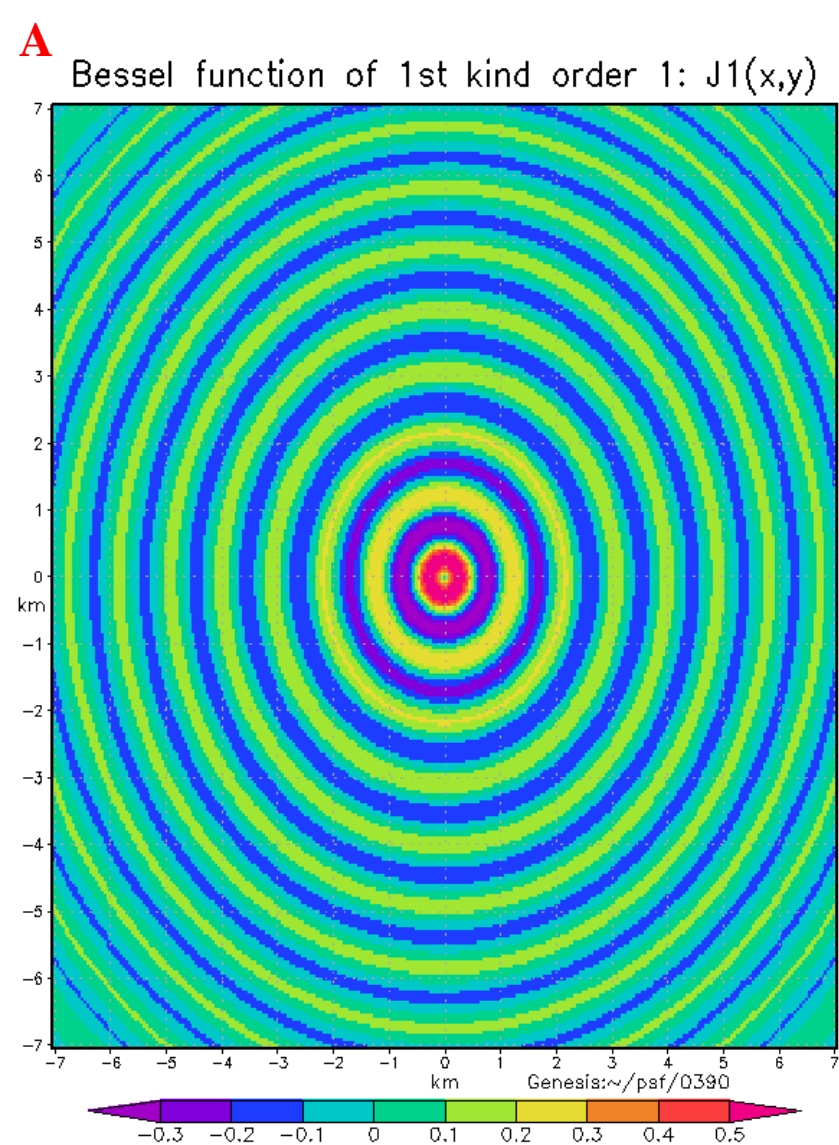


Figure 1: Plots of the Bessel function of the first kind, order 1, J_1 in (A) two dimensions and (B) one dimension. Both the x-axis and y-axis are in km.

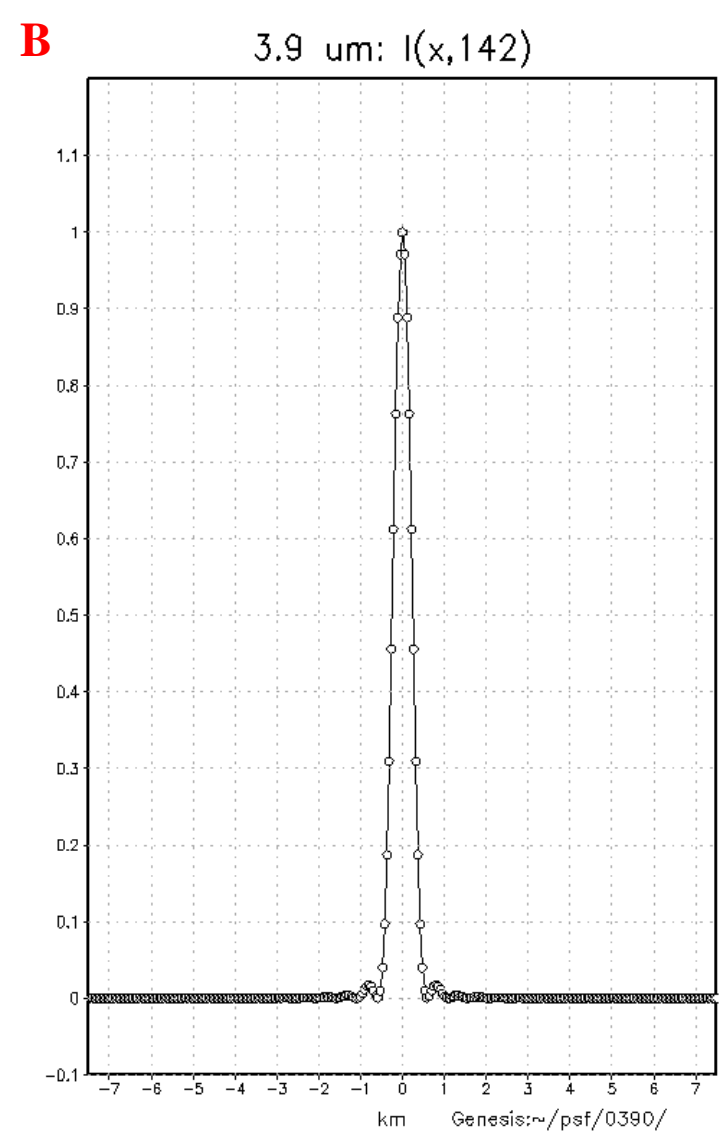
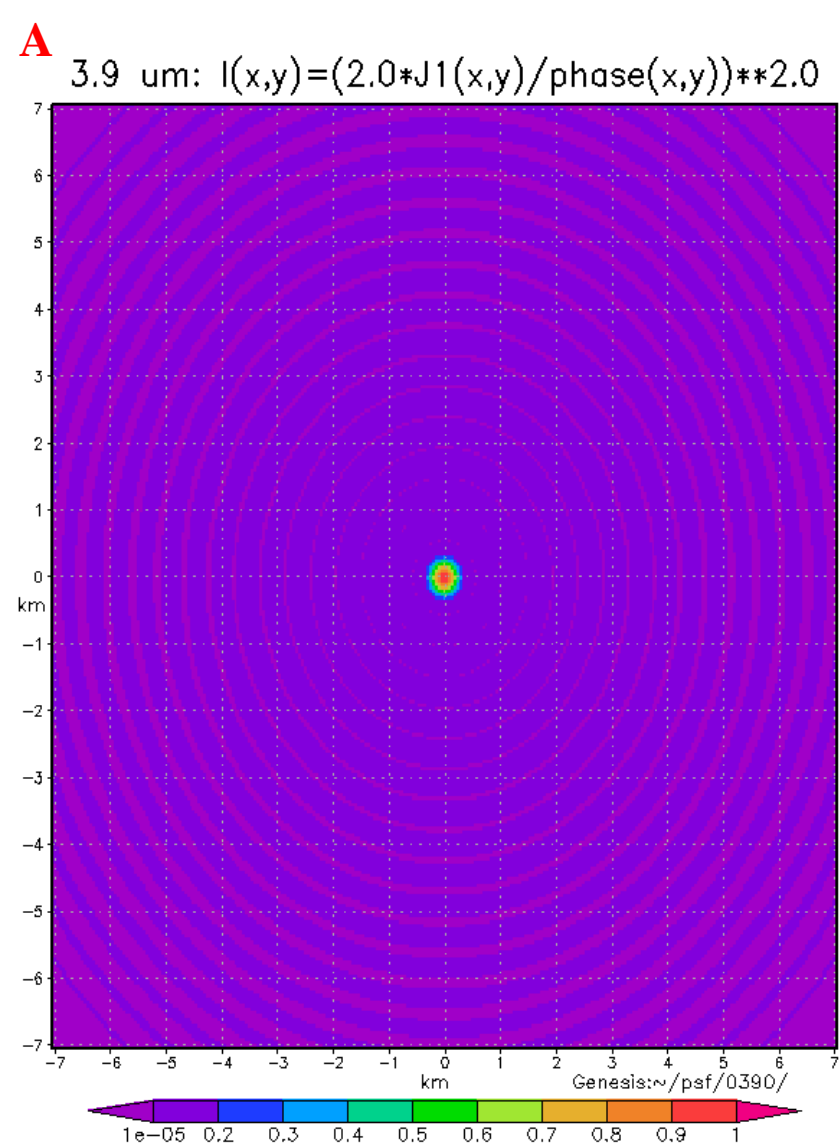


Figure 2: Plots of $I(\theta) = I(0) * (2J_1(\gamma) / \gamma) ** 2.0$, where $\gamma = kR \sin \theta$ and $I(0) = 1.0$ projected on the surface of the Earth, in (A) two dimensions and (B) one dimension. Both the x-axis and y-axis are in km. These plots would show the diffraction pattern on detectors in a satellite by replacing h with f . I think the x-axis would be on the order of the wavelength of the incident radiation: $\sim 1.0 \times 10^{-6}$ m. See slide 4.

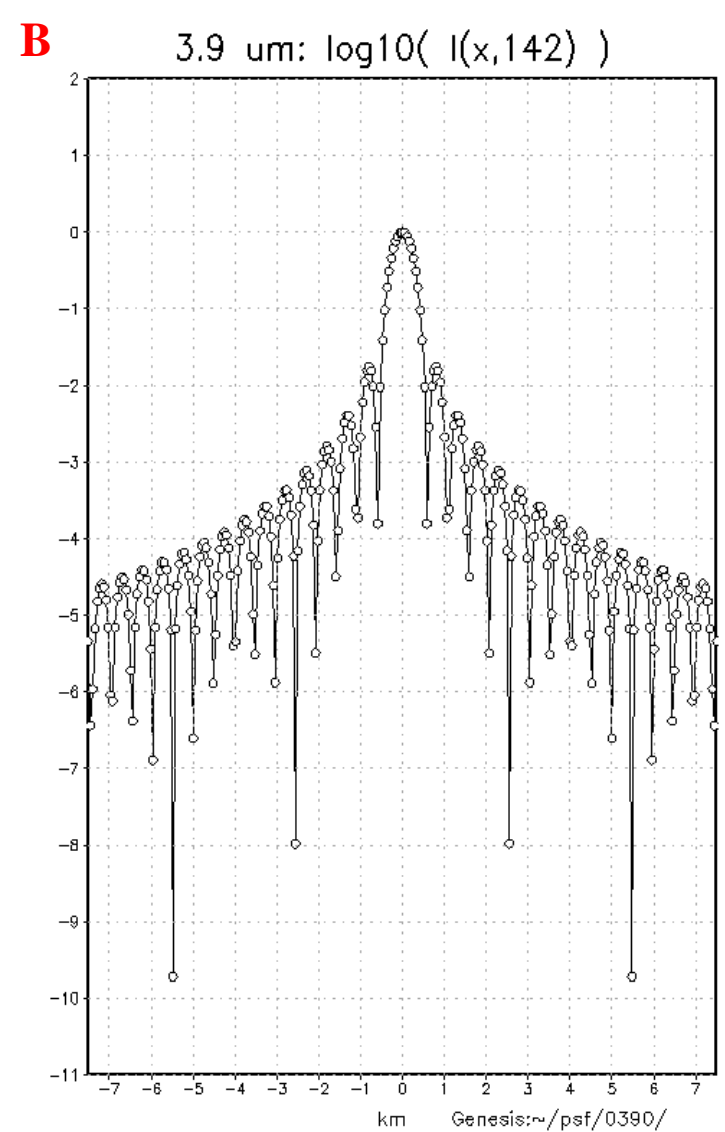
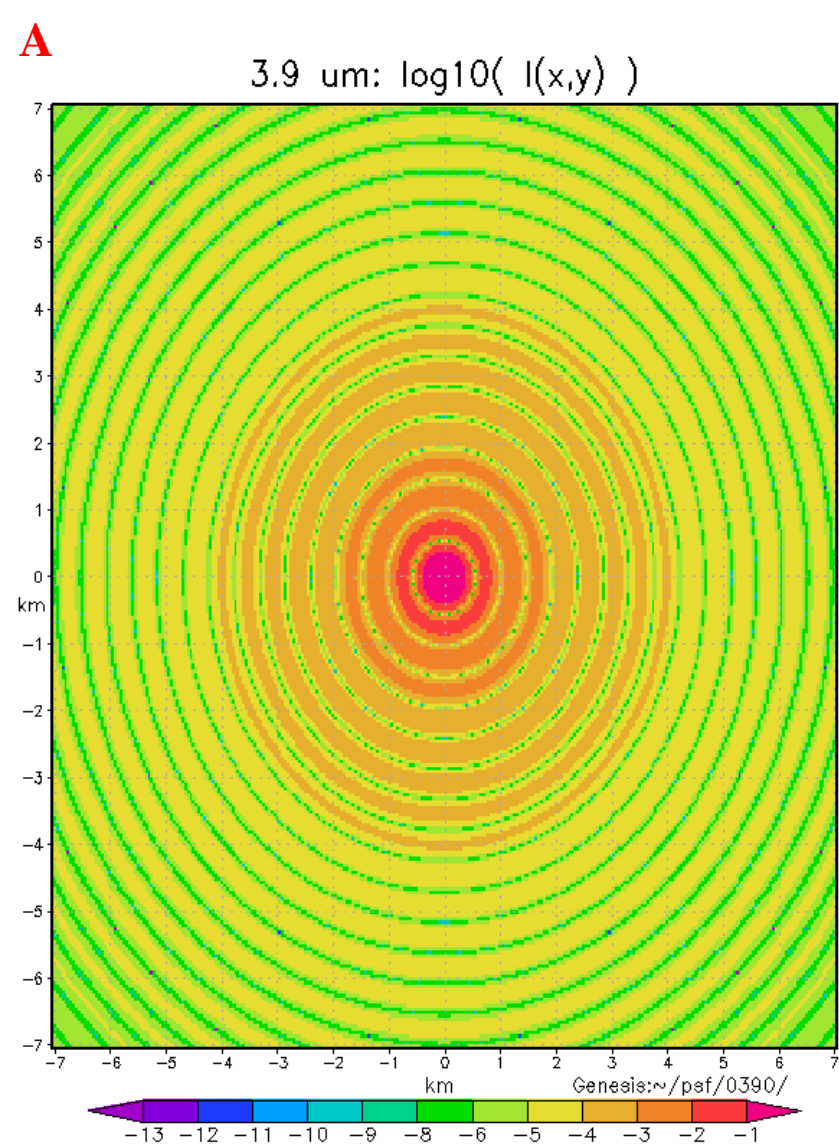


Figure 3: Plots of $\log_{10}(I(\theta))$ (A) two dimensions and (B) one dimension. Both the x-axis and y-axis are in km. Downward “spikes” in (B) represent undefined values of $\log_{10}(I(x,y)=0)$.

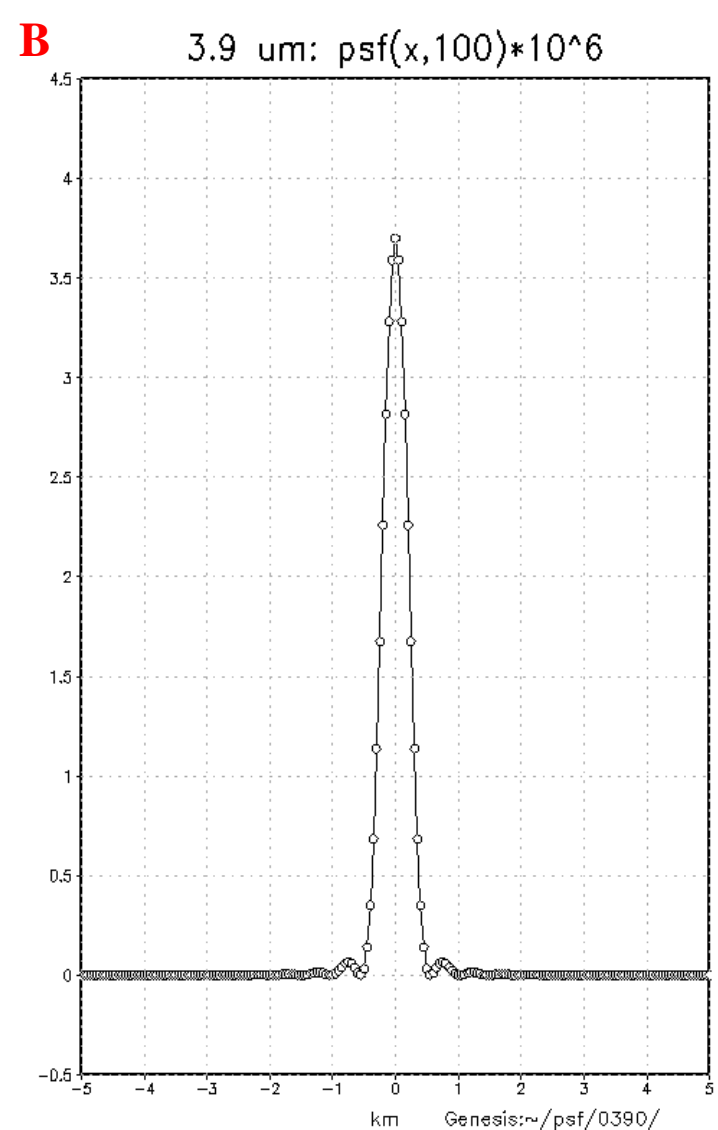
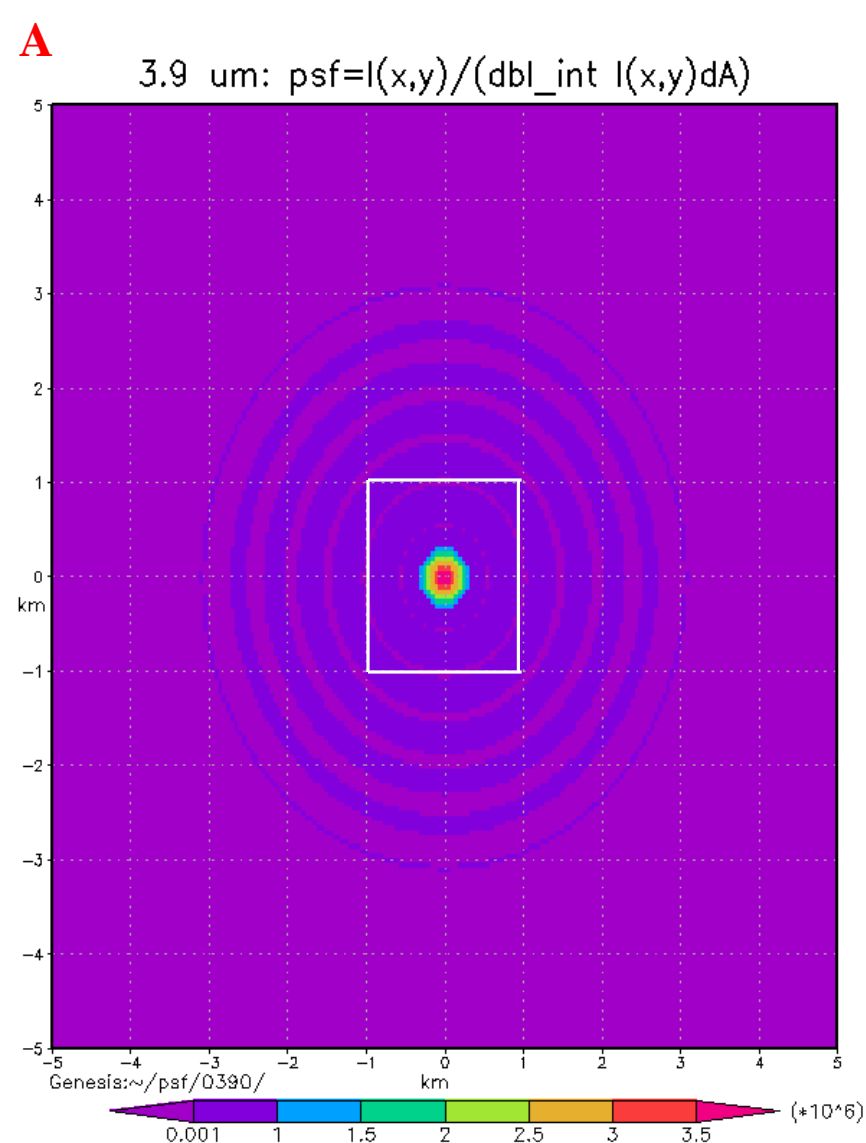


Figure 4: Plots of PSF (**A**) two dimensions and (**B**) one dimension. Both the x-axis and y-axis are in km. Red box in (**A**) is one GOES-16 footprint. **These plots show PSF(x,y) on the surface of the Earth.** Note: PSF(x,y) here (Zhang et al. 2006) excludes the Δt measurement of energy during satellite scanning. Zhang et al. Impact of PSF on Infrared Radiances from Geostationary Satellites, IEEE Trans.Geosci. Remote Sens., vol. 44, no. 8, August 2006

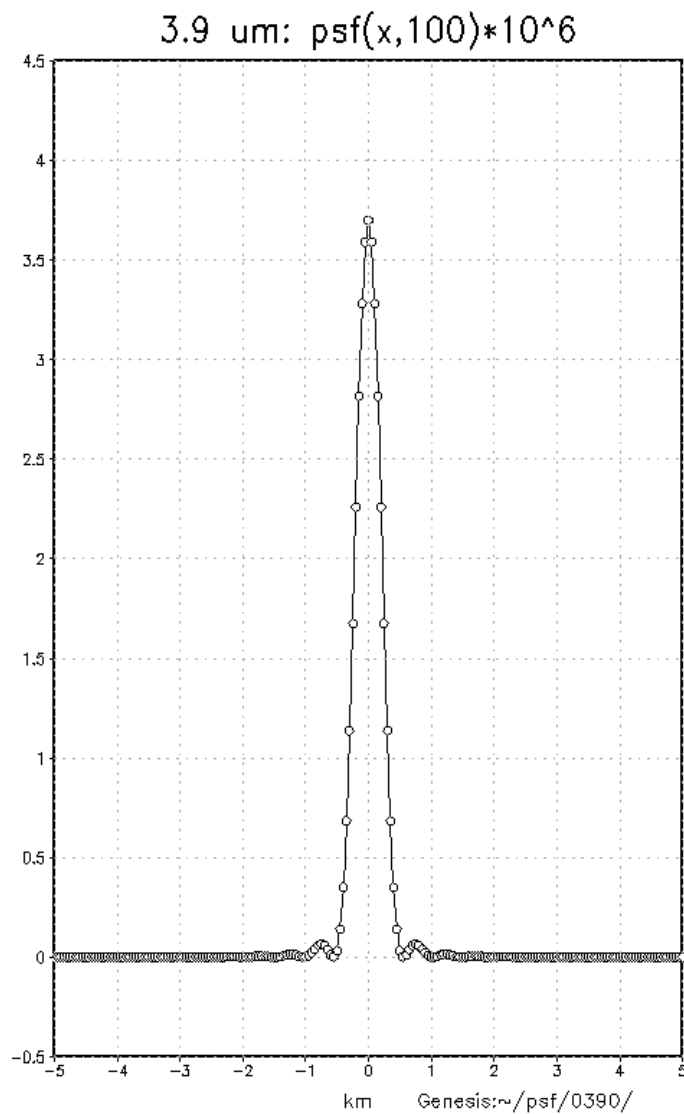
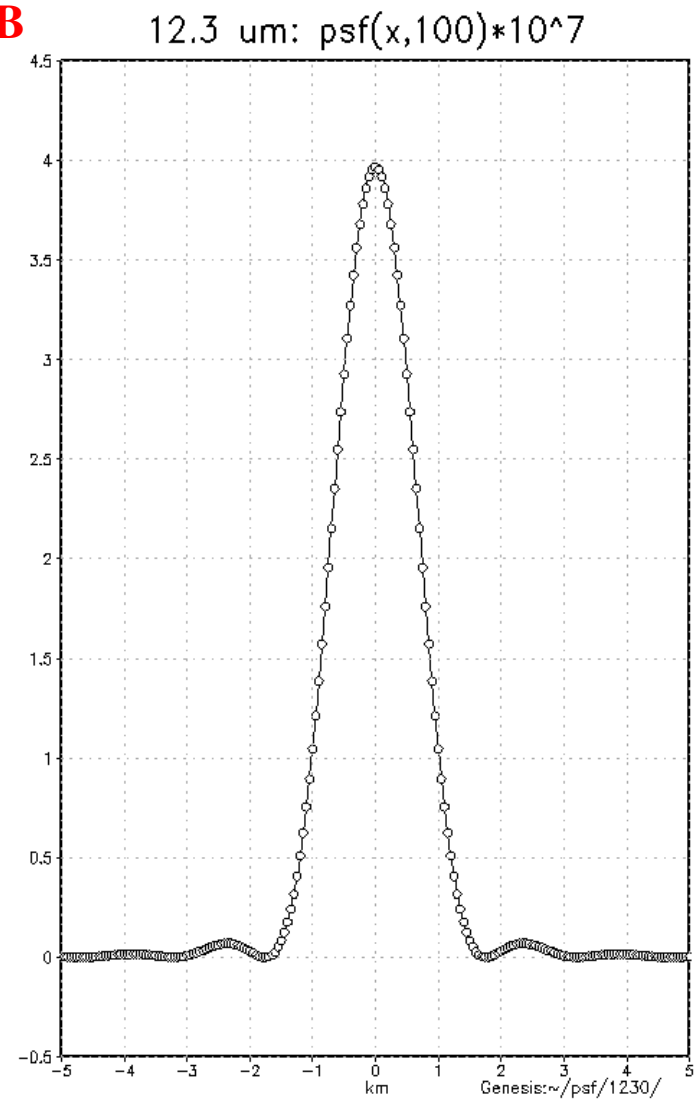
A**B**

Figure 5: A comparison of PSFs at 3.9 μm and 12.3 μm . From slide 4: $\sin\theta = \lambda\gamma/\pi d$. Thus, the width of the Airy disk is proportional to wavelength. This is why the Airy disk at 12.3 μm is larger than that at 3.9 μm . Note the Airy disk is larger than the ABI 2 km footprint. **Alternately,** the intensity at 3.9 μm /12.3 μm is concentrated/spread out on a focal plane of a satellite.

Diffraction patterns for VIIRS M12 (3.7 μ m)

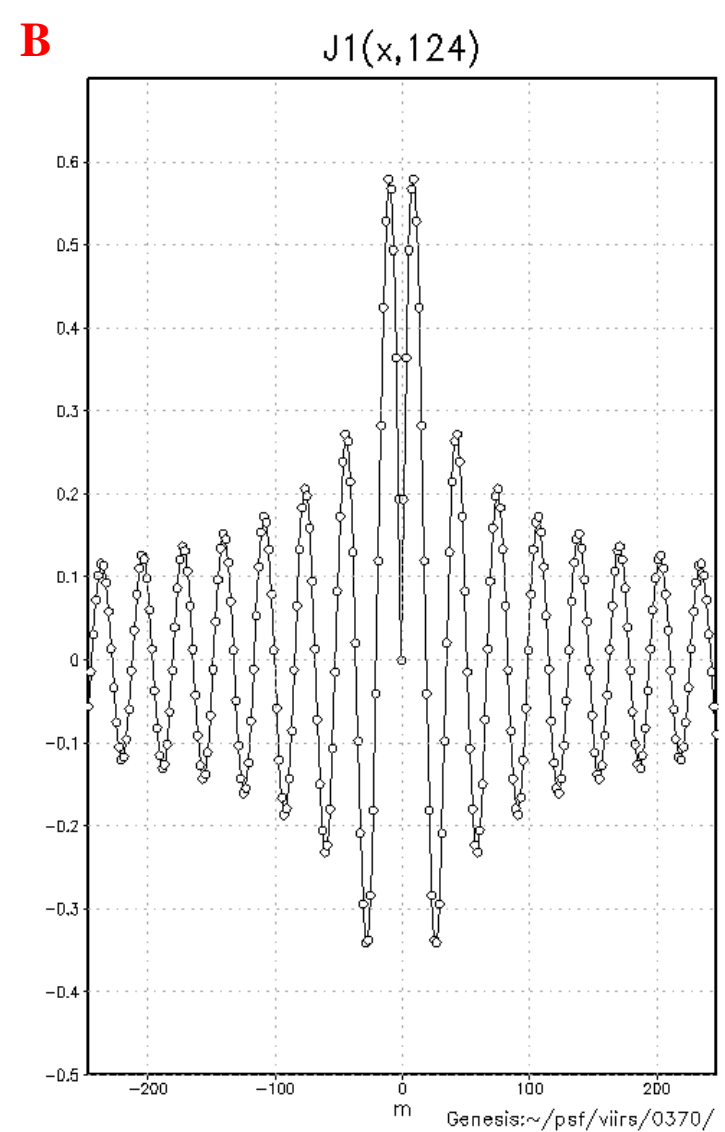
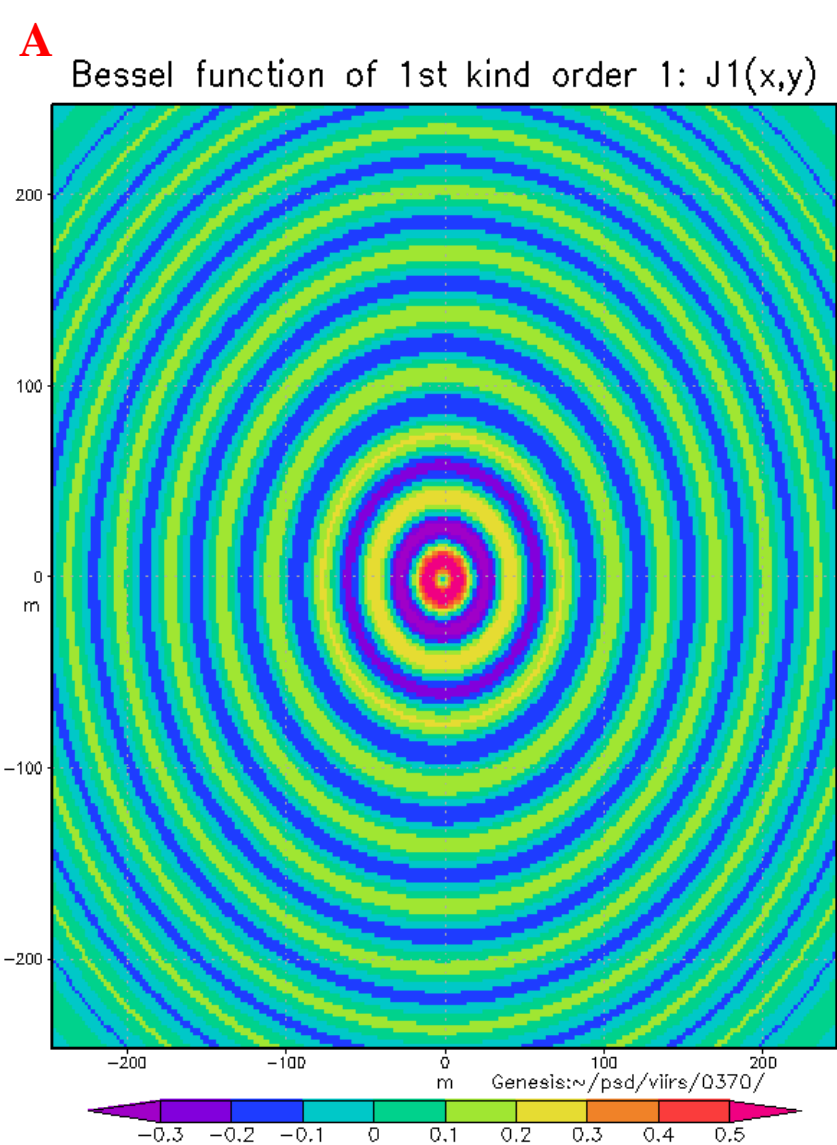


Figure 6: Plots of the Bessel function of the first kind, order 1, J_1 in (A) two dimensions and (B) one dimension. Both the x-axis and y-axis are in **meters** for VIIRS M-Bands: *Footprint~750 m*.

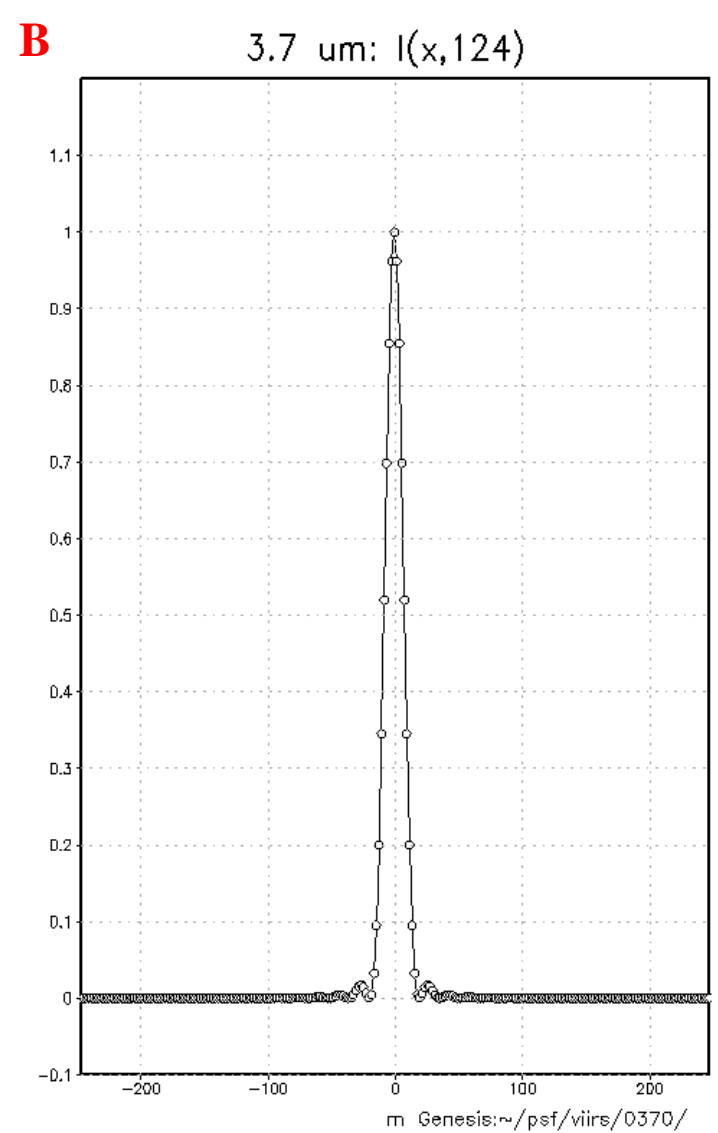
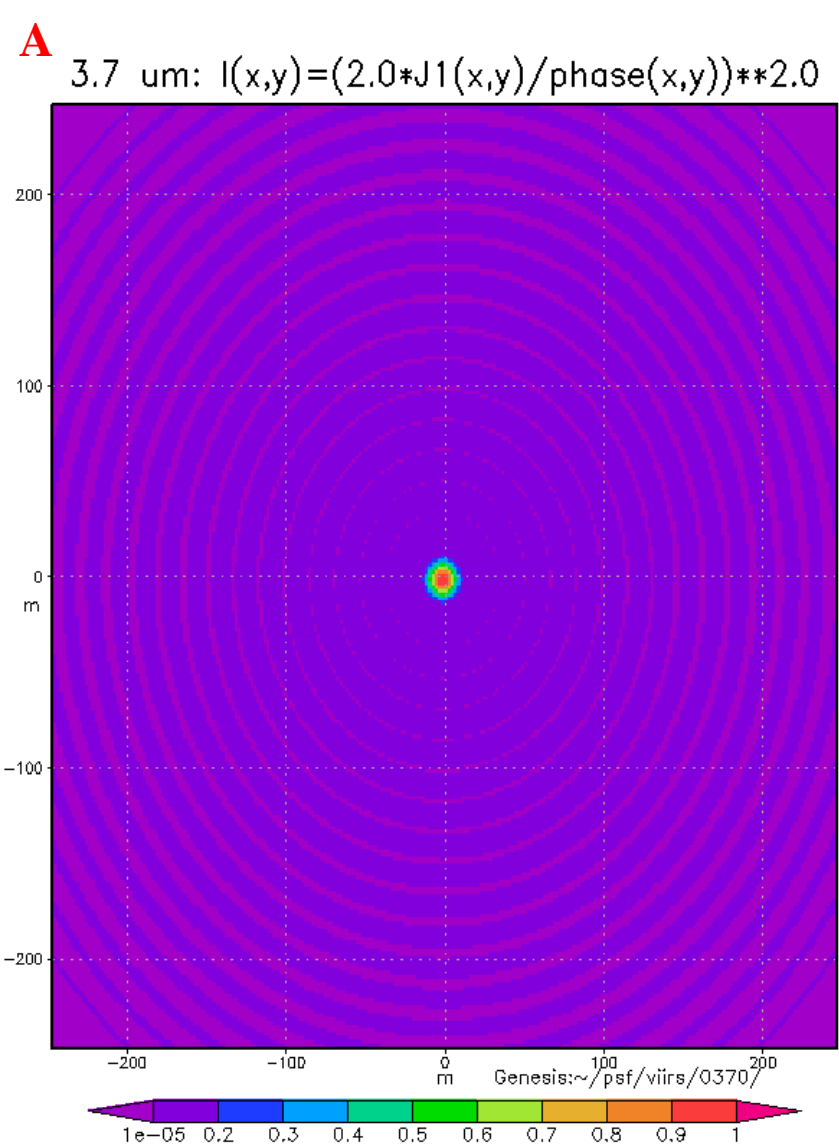


Figure 7: Plots of $I(\theta) = I(0) * (2J_1(\gamma) / \gamma) ** 2.0$, where $\gamma = kR \sin \theta$ and $I(0) = 1.0$ projected on the surface of the Earth, in **(A)** two dimensions and **(B)** one dimension. Both the x-axis and y-axis are in **meters** for VIIRS M-Bands: *Footprint*~750 m.

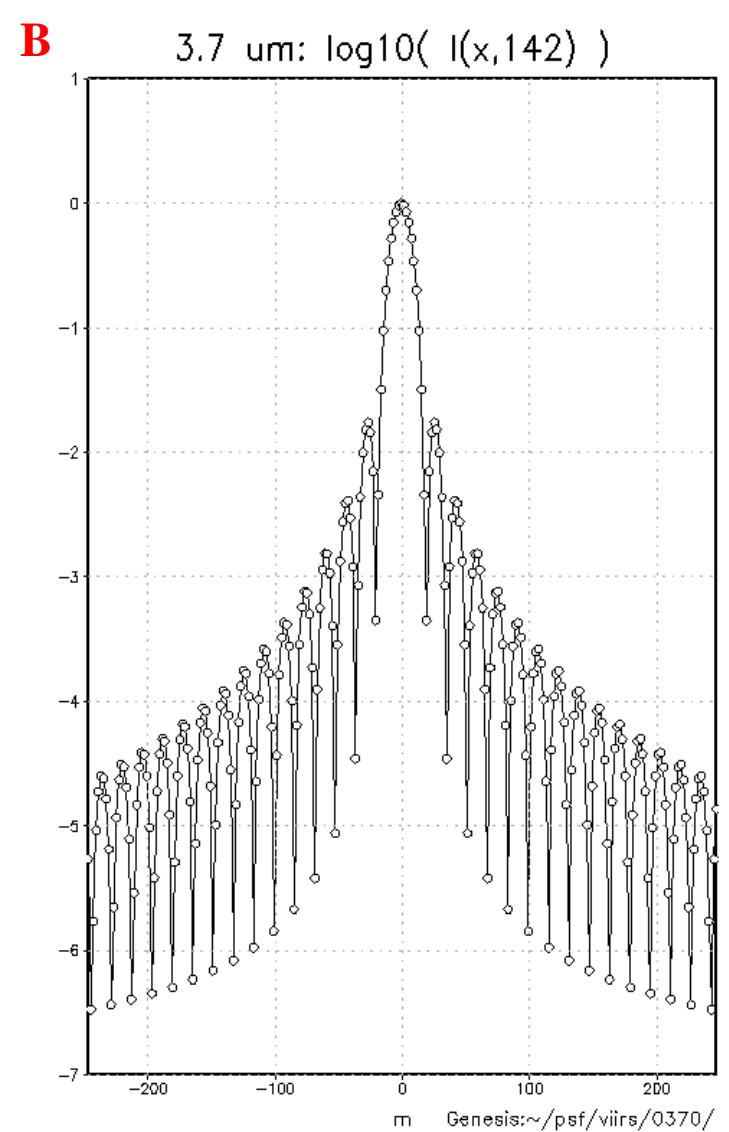
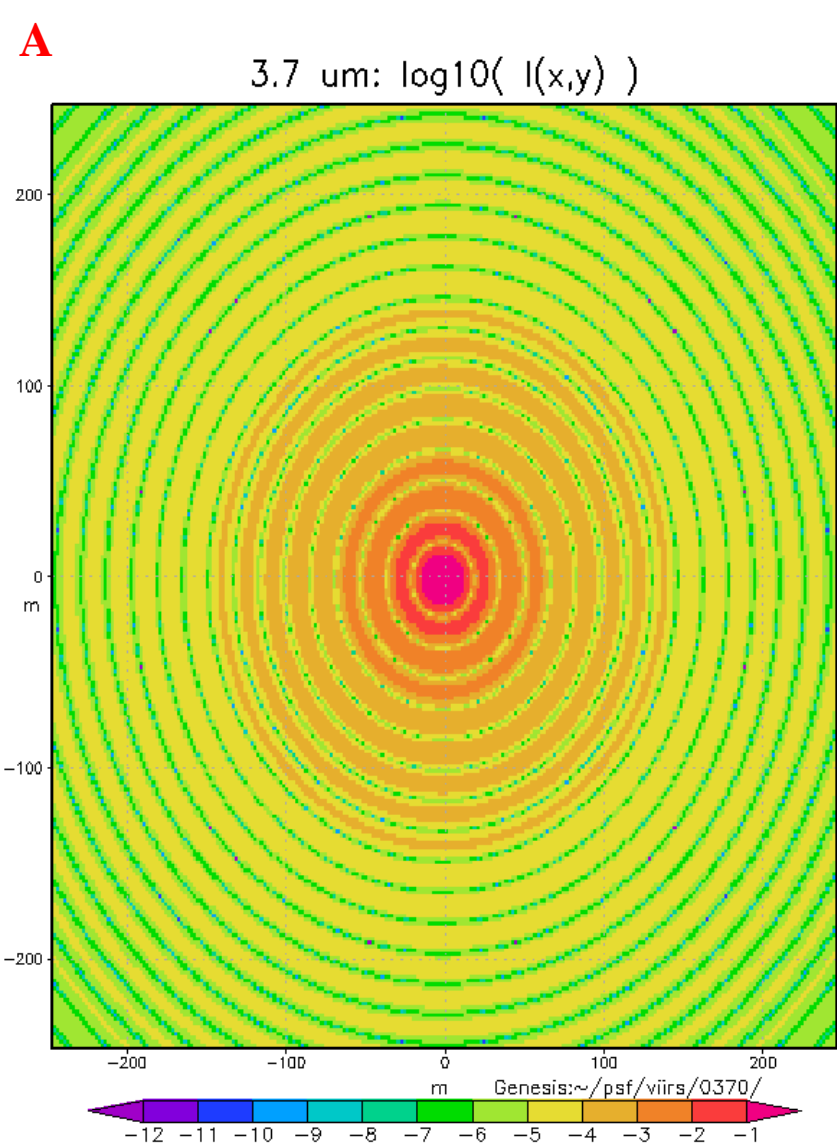


Figure 8: Plots of $\log_{10}(I(\theta))$ (**A**) two dimensions and (**B**) one dimension. Both the x-axis and y-axis are in **meters** for VIIRS M-Bands: *Footprint*~750 m. Downward “spikes” in (**B**) represent undefined values of $\log_{10}(I(x,y)=0)$.

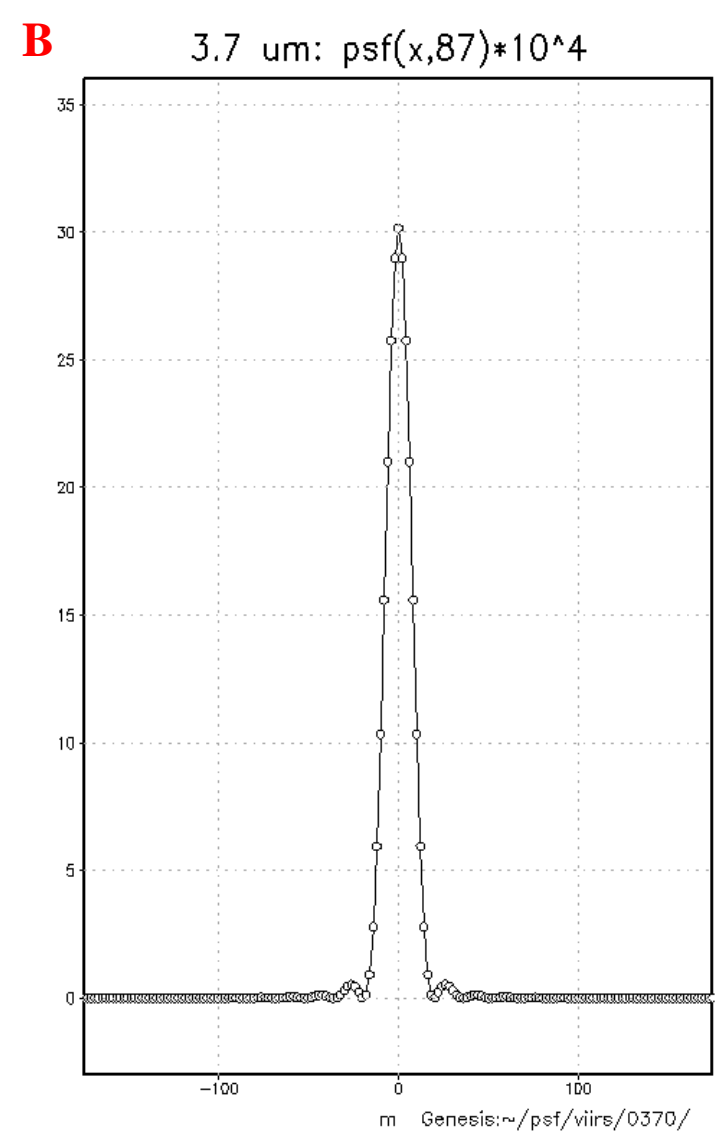
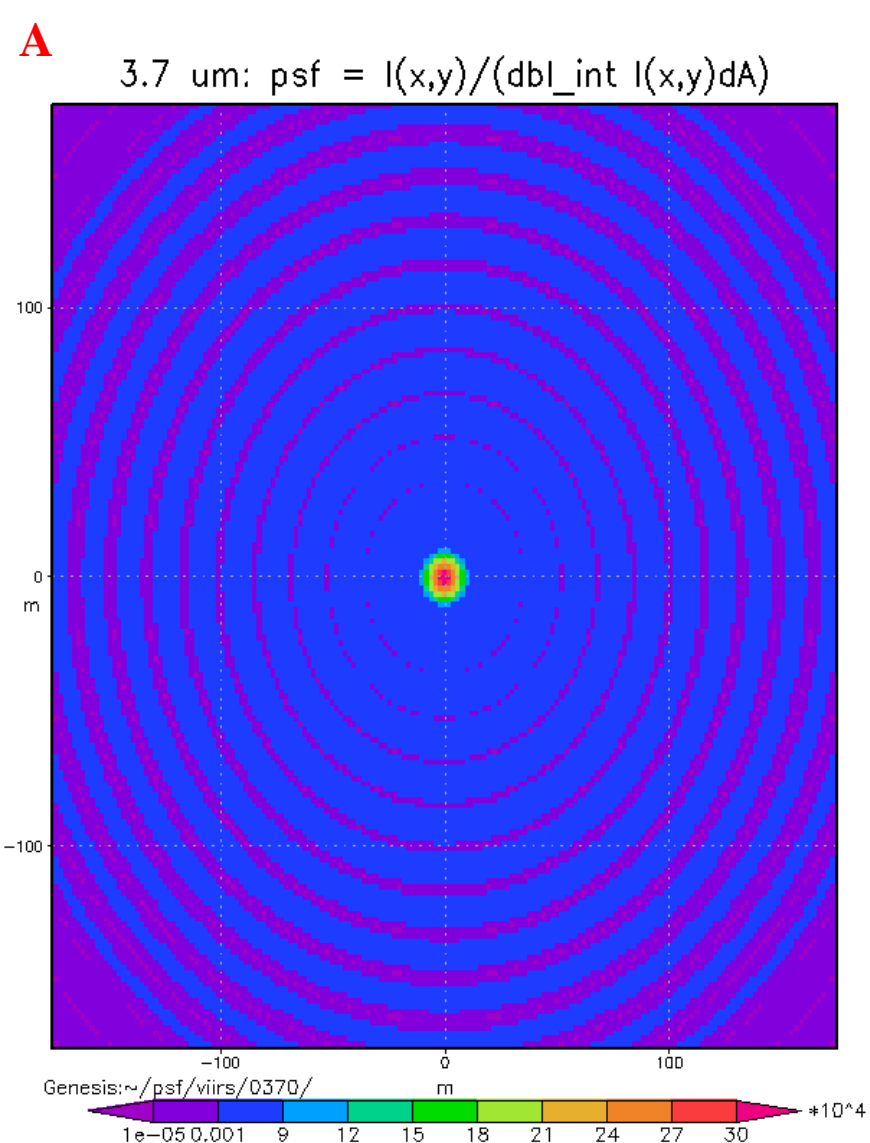


Figure 9: Plots of PSF (**A**) two dimensions and (**B**) one dimension. Both the x-axis and y-axis are in **meters** for VIIRS M-Bands: *Footprint*~750 m. **These plots show PSF(x,y) on the surface of the Earth.** Note: PSF(x,y) here (Zhang et al. 2006) excludes the Δt measurement of energy during satellite scanning. Zhang et al. Impact of PSF on Infrared Radiances from Geostationary Satellites, IEEE Trans.Geosci. Remote Sens., vol. 44, no. 8, August 2006

From slide 7, Fig.2 caption, we have $\gamma = kR \sin\theta = (\pi D/\lambda) * \sin\theta$. Solve for $\sin\theta$ yields, $\sin\theta = \gamma\lambda/\pi D$. Using the small angle approximation, $\sin\theta \sim \theta$. We now have $\theta = \gamma\lambda/\pi D$. From the figure on slide 2, $\theta = d/h$. Since $\theta = \theta$, we have $d/h = \gamma\lambda/\pi D$. Solving for d we get

$$d = \gamma\lambda h/\pi D.$$

$\gamma = 3.832$ gives the first zero of $J_1(\gamma)$.

$\lambda = 3.9\text{e-}6 \text{ m}$ (ABI), $= 3.7\text{e-}6 \text{ m}$ (VIIRS M-Band) is the wavelength of incident radiation.

$h = 3.5786\text{e}7 \text{ m}$ (ABI), $= 8.24\text{e}5 \text{ m}$ (VIIRS) is the height of sensor above the surface of the Earth.

π is very irrational.

$D = 3.048\text{e-}1 \text{ m}$ (ABI), $= 1.91\text{e-}1 \text{ m}$ (VIIRS) is the aperture diameter.

$d = 558 \text{ m}$ (ABI), $= 19.47 \text{ m}$ (VIIRS M-Band) is the **radius** of the Airy disk.

NOTE

The **diameter** of an Airy disk to an ABI footprint is $(2*558/2000)*100.0 = 55.8 \%$.

The **diameter** of an Airy disk to an M-Band footprint is $(2*19.47/750)*100.0 = 5.18 \%$.

VIIRS has a focal length of $f = 1.14 \text{ m}$. We can use similar triangles (Slide 2) to get, $d(\text{detector}) = (1.14*19.47)/824000 = 0.27 \text{ }\mu\text{m}$. Thus the angular measure of the detector would be $\theta = d(\text{detector})/f = 0.27 \text{ }\mu\text{m}/1.14\text{m} \sim 0.24 \text{ }\mu\text{rad}$. We can also compute $\theta = d/h = 19.47\text{m}/82400\text{m} \sim 0.24 \text{ }\mu\text{rad}$.

Thus, detector **radius** $= 0.27 \text{ }\mu\text{m}$ or $\theta(\text{half width}) = 0.24 \text{ }\mu\text{rad}$ for VIIRS: Can anyone confirm this result?

My Thoughts

- A diffraction pattern of the intensity is what the sensors, in the optics of a satellite, measure. When the measurement occurs for a time Δt , a $PSF(x,y,t)$ results for scanning (pick your broom type) sensors.
- Some refer to the PSF as being “on the surface” of the Earth and neglect Δt sampling. For example,
- Zhang et al. Impact of PSF on Infrared Radiances from Geostationary Satellites, IEEE Trans.Geosci. Remote Sens., vol. 44, no. 8, August 2006
- Grandell, J. and R. Stuhlmann, 2007: Limitations to a Geostationary Infrared Sounder due to Diffraction: The Meteosat Third Generation Infrared Sounder (MTG IRS). J. Atmos. Ocean Tech., vol 24, pp 1740-1749.
- Huang, C, J. R.G. Townshend, S. Liang, S. N.V. Kalluir, and R. S. DeFries, 2002: Impact of sensor’s point spread function on land cover characterization: Assessment and deconvolution. Remote Sensing of Environment, vol 80, pp 203-212.
- Fires outside an ABI central footprint may influence the radiance of the central footprint via the $PSF(x,y,t)$. Since the *diameter* of the Airy for a VIIRS M-band is so small (5.18%) relative to an M-band footprint, can fires outside a central VIIRS footprint influence the radiance of a central footprint?

Applications

- 1) Build $\text{PSF}(x,y,t)$ for any satellite and any wavelength. Need height of satellite from the surface of the Earth, wavelength of incident light, telescope aperture, satellite image footprint size and dx of PSF. You pick dx .
- 2) Apply $\text{PSF}(x,y,t)$ at $\sim 3.9 \mu\text{m}$ to fires of, which have a size that is a multiple of dx meters. This application will be used to examine the sensitivity of fire hotspots, “under” the $\text{PSF}(x,y,t)$, to build one JPSS footprint. Since the size of the Airy disk is so small relative to one VIIRS M-band footprint (5.18 %, slide 16), can we actually use the “surface of the Earth” $\text{PSF}(x,y,t)$ concept? Do we need to know more about how the sensor (detector) works?
- 3) Ed Hyer (NRL, Monterey) just visited and delivered a presentation in which the above application has already been done by

Schroeder, W., I. Csiszar, L. Giglio, and C. C. Schmidt (2010), On the use of fire radiative power, area, and temperature estimates to characterize biomass burning via moderate to coarse spatial resolution remote sensing data in the Brazilian Amazon, *J. Geophys. Res.*, 115, D21121, doi:10.1029/2009JD013769.

Thank You

